the same temperature, while there is an increase of more than a factor of two at $274^{\circ}\kappa$ compared to $80^{\circ}\kappa$ in TlBr and TlCl. Furthermore, there is also good agreement within experimental error between the values obtained using single crystals and evaporated films.

Fig. 3

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Absorption curves of the E_0 exciton in TlBr at various hydrostatic pressures and at 80° κ . a, 0.65; b, 2.07; c, 3.45; d, 4.85 kbars.

Table 2.	The hydrostatic pressure coefficients for exciton absorption peaks
	in the thallous halides

Material film (f) or crystal (c)	Peak position (ev)	$\stackrel{\rm Temperature}{(^{\circ}\kappa)}$	Pressure coefficient ($\times 10^{-6} \text{ ev/bar}$)
TICI f	$E_{0} 3.40$	80	-8.6 + 0.4
f	E ₀ 3.50	274	$-22 \cdot 2 \div 1 \cdot 0$
TlBr c	$E_{0} 3.030$	80	-9.5 ± 0.4
f	$E_{0} 3.030$	80	-9.2 ± 0.4
f	$E_0 3.065$	195	-11.3 ± 0.5
e	$E_0 3.097$	274	-20.1 ± 1.0
c	$E_{1} 4.064$	80	-0.6 ± 0.4
f	E_{2}^{*} 5.06	80	$-4\cdot 1 \pm 0\cdot 7$
TlI f	$E_{0} 2.80$	80	-8.6 ± 0.4
f	$E_{0}' 3.50$	80	-8.1+0.5

3.3. Pressure Coefficients for Higher Energy Transitions above the Band Edge

The pressure coefficients at 80° K for the absorption peaks E_1 and E_2 in TlBr are given in table 2. The variation of these peak energies, together with E_0 , has been plotted versus pressure and lattice spacing in fig. 4. The lattice spacing has been calculated assuming a volume isothermal compressibility which is constant with respect to pressure : $\beta_{T=80^{\circ}\text{K}} = 3.89 \times 10^{-6}$ /bar[†]. The peak E_1 was found to be considerably less sensitive to pressure than E_0 and E_2 . Together with the temperature dependence of these absorption peaks, this supports the view that the E_1 and $E_1 + \Delta$ transitions involve interband states of different symmetry from the remaining strong transitions.





The variation with pressure and lattice spacing of the energy of exciton absorption peaks in single crystals of TlBr an $80^{\circ}\kappa$.

The next higher energy transition (E_0') after the minimum gap in TII has also been studied under pressure. Table 2 shows that the pressure coefficients of E_0 and E_0' at 80° κ are similar.

[†] This value for the isothermal compressibility has been calculated from the adiabatic compressibility and density data of Vallin, Marklund and Sikström (1966); the specific heat per mole at constant pressure quoted by Kelley (1934); and the usual relation between adiabatic and isothermal compressibility (Morse and Lawson 1967). The specific heat and the expansion coefficient are assumed to obey the Grüneisen formula (Ziman 1964).